

EM waves

Polarization conventions

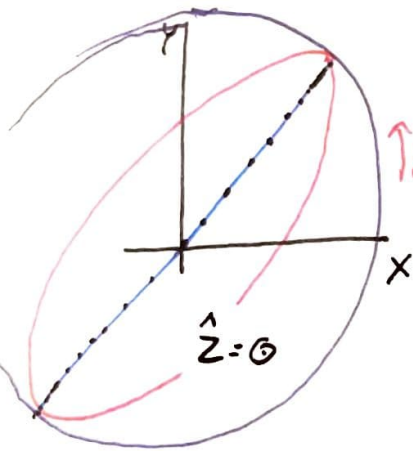
Definition of Terms

E field plane wave can be written

$$\vec{E} = (E_x \hat{x} e^{i\phi_x} + E_y \hat{y} e^{i\phi_y}) \exp(i(kz - \omega t))$$

at fixed z (or fixed t) $\vec{E}$ traces out an ellipse.

Relative phase $\delta \equiv \phi_x - \phi_y$ ^{+E vector sets} the shape

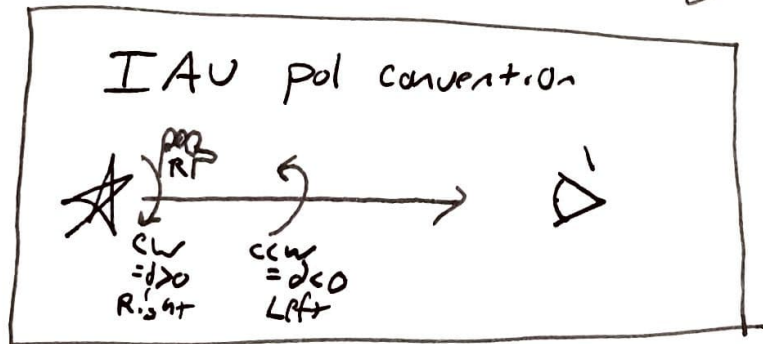


$\phi_x = \phi_y = 45^\circ$ Line

$\phi_x = 0, \phi_y = 90^\circ$ circle

clockwise from observer

$E_x = E_y$



In general

E field amp also matters

polarization usually measured in Stokes parameters

$$I = \langle E_x^2 + E_y^2 \rangle / R_0$$

$$Q = \langle E_x^2 - E_y^2 \rangle / R_0$$

$$U = \langle 2E_x E_y \cos \delta \rangle / R_0$$

$$V = \langle 2E_x E_y \sin \delta \rangle / R_0$$

$$R_0 = \frac{4\pi}{c}$$

• Avg over 1 cycle
 $\tau \gg (1/\omega)$
 • assuming stochastic source

EM cont

Def: polarized flux

$$I_p = \sqrt{(Q^2 + U^2 + V^2)}$$

polarization fraction

$$p = \frac{I_p}{I}$$

Note that Stokes parameters are also definable
in terms of circular polarizations. •

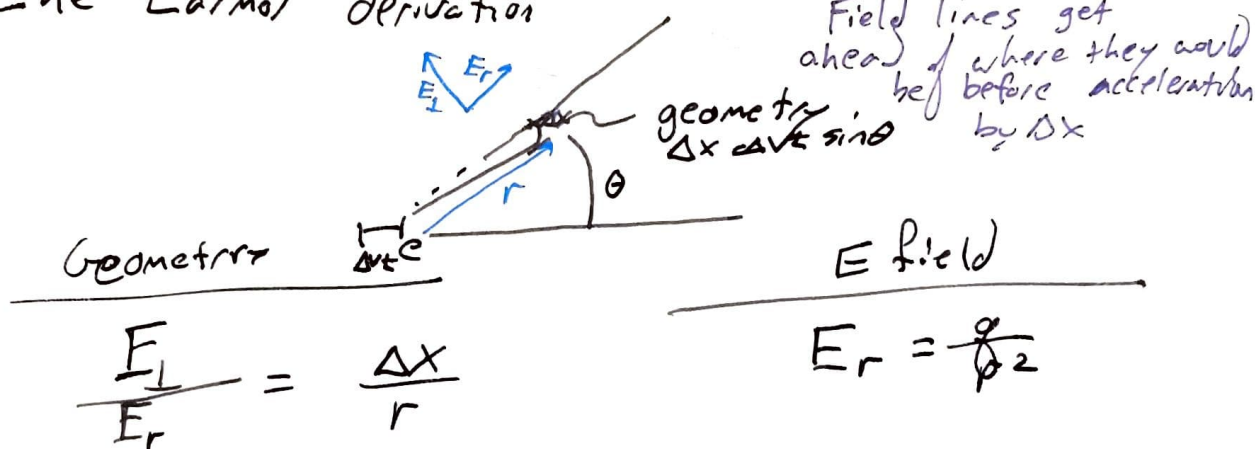
Antennas

ERA 2.7
Condon + Pansan

Moving e^-
Accelerating e^- s generate propagating E fields



The Larmor derivation



$$t = \frac{r}{c}$$

+

$$= \frac{\Delta v \sin \theta}{c \Delta t}$$

$$\frac{\Delta v}{\Delta t} \rightarrow \dot{v}$$

$$E_{\perp} = \frac{q \dot{v} \sin \theta}{r c^2}$$

power pattern $\rightarrow$ Poynting flux

$$\vec{S} = \frac{c}{4\pi} \vec{E} \times \vec{B}$$

in CGS $|\vec{E}| = |\vec{B}|$

$$S = \left(\frac{q^2 \dot{v}^2}{4\pi c^3} \right) \frac{\sin^2 \theta}{r^2}$$

Integrate over sphere

... (two r^2 in dA cancel)



$$P = \frac{2}{3} \frac{q^2 \dot{v}^2}{c^3}$$

Note, do not use if

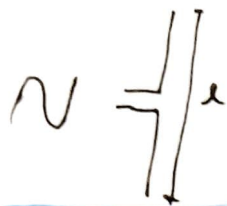
$v \sim c$

$\Delta r \sim$ bohr radius

More Antennas

ERA 3.1

A dipole



Drive the dipole w/ a sinusoid

$$I = I_0 e^{-i\omega t}$$

The moving charges radiate according to Larmor.

dz an infinitesimal charge moves

$$dE_{\perp} = \frac{dq \dot{v} \sin \theta}{rc^2}$$

If $l \ll \lambda$



Then all dz charges are moving in phase. $\therefore E_{\perp}$ s will add up

$$E_{\perp} = \int_{-l/2}^{l/2} \frac{\dot{v} \sin \theta}{rc^2} dq$$

$$= \int_{-l/2}^{l/2} \frac{dq}{dz} dz \frac{\dot{v} \sin \theta}{rc^2}$$

@ $r \gg l$, the 'far' field r can be same for all dq .

use our sinusoid with $\dot{v} = -i\omega v$

$$E_{\perp} = \frac{-i\omega \sin \theta}{rc^2} \int \frac{dq}{dz} v dz$$

$$\frac{dq}{dz} \frac{dz}{dt} = I$$

$$E_{\perp} = \frac{-i\omega \sin \theta}{rc^2} I dz$$

Current is going in, not coming out



Must go to zero @ end!

Ants x3



Carson + Renshaw say
"current" declines almost linearly" bc it is the
"tail end of a standing wave sinusoid"

Maybe they mean $\sin(kz)$
where $k = \frac{2\pi}{\lambda}$ if $z \ll \lambda$ then

$$\sin(kz) \approx (1 + \frac{kz}{2})$$

But not quite... anyway, I'm ~~to~~ Scott
Let

$$I(z) \approx I_0 e^{-i\omega t} \left[1 - \frac{|z|}{\lambda/2} \right]$$

$$E_{\perp} = \frac{-i\omega \sin\theta}{rc^2} \int dz$$

$$\approx \frac{I_0 l}{2} e^{-i\omega t} \left(\frac{-i\omega \sin\theta}{rc^2} \right)$$

Whatever bus ~~from~~

Power
(time avg)

$$\langle S \rangle = \frac{c}{4\pi} \left(\frac{I_0 l}{\lambda} \frac{\pi}{c} \right)^2 \frac{\sin^2\theta}{r^2} \langle e^{i\omega t} \rangle$$

~~or~~

$$P \propto \sin^2\theta$$

like the e^- in Larmor
bc we made assumptions $l \ll \lambda$, $e^{i\omega t}$
to make it so. But still. Dipole!

this avg
to $\frac{1}{2}$ bc
we really meant
 $\sin(\omega t)$
the whole time
 $\langle \sin^2(\omega t) \rangle = \frac{1}{2}$

Antennas

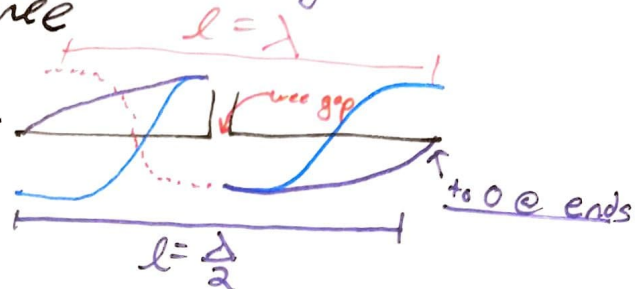
The total power. Int over θ sphere

$$\langle P \rangle = \frac{\pi^2}{3c} \left(\frac{I_0 l}{A} \right)^2$$

length of dipole is wavelengths

Now Dipoles @ Resonance

$$l \approx \frac{\lambda}{2}$$

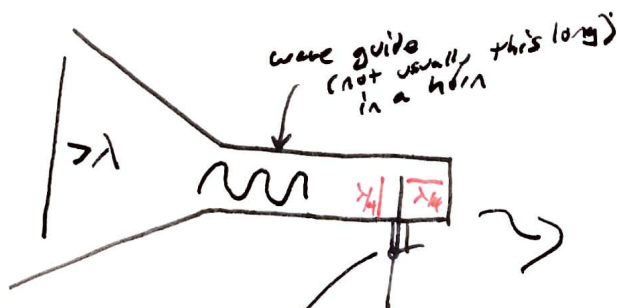
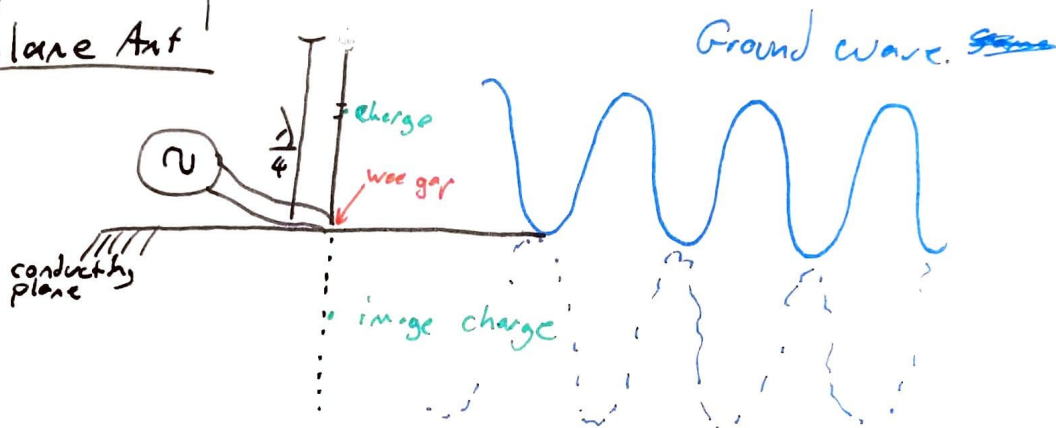


A magic dipole The Groundplane Ant

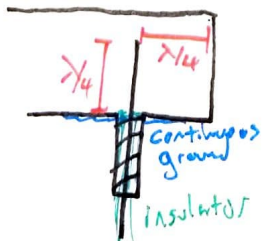
Same pattern as dipole.



Field $E=0$ below gnd.
What happens here??
probably best we don't ask.



Note Feed detail



see ERA 3.4 on waveguides



to keep out horizontal waves.